

LNF - 66/23
13 Aprile 1966

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PHOTOPRODUCTION OF π^+ BY LINEARLY POLARIZED γ
RAYS IN THE ENERGY RANGE 200 - 400 MeV. -

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P. Gorenstein, M. Grilli, M. Nigro^(x), E. Schiavuta^(x), F. Soso, P. Spillantini, V. Valente^(x): ASYMMETRY RATIOS IN THE PHOTOPRODUCTION OF π^+ BY LINEARLY POLARIZED γ RAYS IN THE ENERGY RANGE 200 - 400 MeV.

(Submitted for publication to Physics Letters)

In the framework of a systematic study of the reaction:

$$(1) \quad \gamma + p = n + \pi^+$$

around the first resonance, using linearly polarized photons, we are measuring the asymmetry ratio

$$(2) \quad A(\theta) = \frac{\sigma_{\perp}(\theta) - \sigma_{\parallel}(\theta)}{\sigma_{\perp}(\theta) + \sigma_{\parallel}(\theta)}$$

with the apparatus and technique previously described⁽¹⁾, at different angles and different incident photon energies.

The quantities σ_{\perp} (σ_{\parallel}), in (2), are defined as the differential cross section for the process (1) by photons with the electric vector perpendicular (parallel) to the production plane. The angle θ is the produc

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tion angle of the π in the c. m. system.

We report here our results for $A(90^\circ)$ in the energy interval $E_\gamma = 200 - 400$ MeV. A small part of these results have previously been published⁽¹⁾.

A measurement of $A(90^\circ)$ as a function of E_γ determine the term $I_0(90^\circ)$ of the following expression, as a function of energy:

$$(3) \quad I(\theta) = - \frac{1}{\text{sen}^2 \theta} \frac{K}{q} \frac{\sigma_\perp(\theta) - \sigma_\parallel(\theta)}{2} = \frac{1}{\text{sen}^2 \theta} A(\theta) \sigma(\theta) =$$

$$= I_0(\theta) + I_1(\theta) \cos \theta$$

where

$K = \gamma$ -ray's momentum in the c. m. system
 $q = \pi$'s momentum in the c. m. system

$$\sigma(\theta) = \frac{\sigma_\perp(\theta) + \sigma_\parallel(\theta)}{2} = \text{differential cross section by unpolarized } \gamma \text{ - rays.}$$

The coefficients I_0 and I_1 are different functions of the amplitude of the various terms contributing to the reaction (1).

In particular,

$$(4) \quad I_0(\theta) = \frac{1}{2} \left\{ |\mathcal{F}_3|^2 + |\mathcal{F}_4|^2 + 2 \text{Re}(\mathcal{F}_1 \mathcal{F}_4^*) + 2 \text{Re}(\mathcal{F}_2 \mathcal{F}_3^*) \right\}, \text{ where}$$

the amplitudes $\mathcal{F}_1, \dots, \mathcal{F}_4$ are defined in (2)

A specific model must be introduced to get definite information from such kind of analysis on I_0 : here we want just to note that the quadratic form of the amplitudes $\mathcal{F}_i (i = 1, \dots, 4)$ entering I_0 is essentially different from the one associated to the differential cross section as given by the following expression:

$$(5) \quad \frac{K}{q} \sigma(90^\circ) =$$

$$= \left[|\mathcal{F}_1|^2 + |\mathcal{F}_2|^2 + \frac{1}{2} |\mathcal{F}_3|^2 + \frac{1}{2} |\mathcal{F}_4|^2 + \text{Re} \mathcal{F}_1 \mathcal{F}_4^* + \text{Re} \mathcal{F}_2 \mathcal{F}_3^* \right]$$

The comparison of the expressions (4) and (5) is a clear example of the complementarity of measurements on process (1) by polarized and unpolarized γ rays, as has been extensively discussed by many authors^(3,6).

The source of polarized photons in our experiment was the coherent bremsstrahlung beam developed at the Frascati synchrotron by Barbiellini et al. ⁽⁷⁾.

The polarization of the beam is defined as

$$(6) \quad P = \frac{N_{\perp} - N_{\parallel}}{N_{\perp} + N_{\parallel}}$$

where N_{\perp} (N_{\parallel}) is the relative number of photons having their electric vector perpendicular (parallel) to the (γ, π) plane. In our measurements the value of P ranged from 33% to 13%. The values of P are not measured directly, but are obtained from Born approximation calculations of coherent bremsstrahlung. The reliability of the values of P lies on the agreement between the experimental and the computed values of the bremsstrahlung intensity. Such calculations are affected by some factors⁽⁸⁾: namely, the incertitude in the knowledge of crystal orientation, the mosaic structure of the crystal, the incoherent background and the atomic form factor. The uncertainties in the calculated values of P due to these factors are, however, certainly smaller than $\Delta P \approx \pm 1\%$.

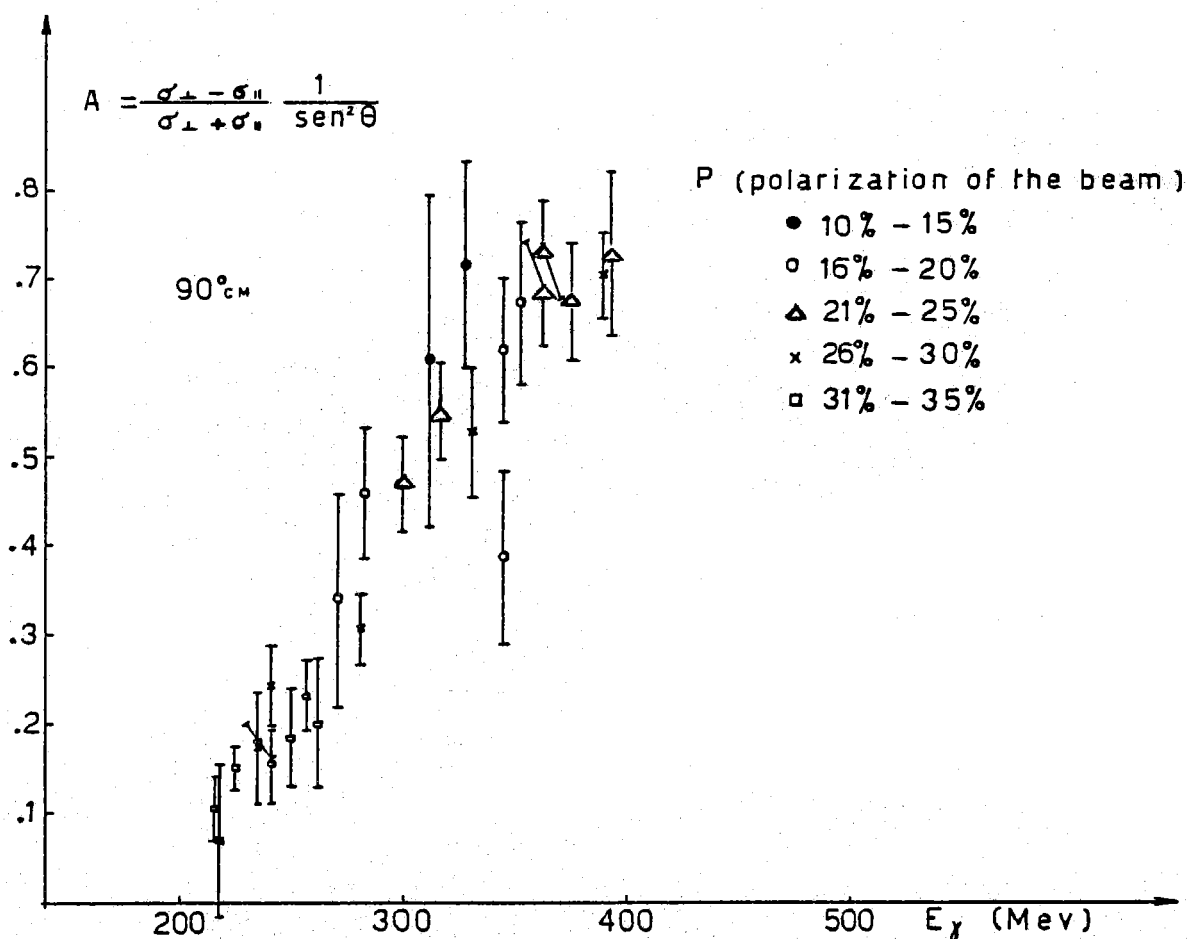


FIG. 1 - All measured points of $A(90^\circ)$ are reported vs γ -ray's energy. Some of the neighboring points have been measured, for comparison, using different values for the polarization of the photon's beam (P).

The data we report for the asymmetry ratio $A(90^\circ)$ have been corrected for the various backgrounds (empty target, electrons, multi-pion photoproduction) as explained in (1). The errors shown are the statistical error only. The total errors, including also the estimated error on P and on the background subtraction, would be at most 1.3 times the statistical errors. Each value of $A(90^\circ)$ has been measured many times in different runs. Moreover some neighboring points have been measured using different values of P . (see fig. 1).

All our measurements are shown in fig. 1. It can be seen that the points taken with different values of P are consistent.

In table I we give our final results for $A(90^\circ)$ at different energies, with points at neighbouring energies lumped together. In the same table we give also previous results of Smith and Mozley⁽⁹⁾.

In fig. 2 all these results for $A(90^\circ)$ are compared with the theoretical prevision of W. Schmidt⁽¹⁰⁾ and of A. Donnachie and G. Shaw⁽³⁾.

It is evident that the results for $E_\gamma > 300$ MeV give a larger asymmetry than is predicted by the theory.

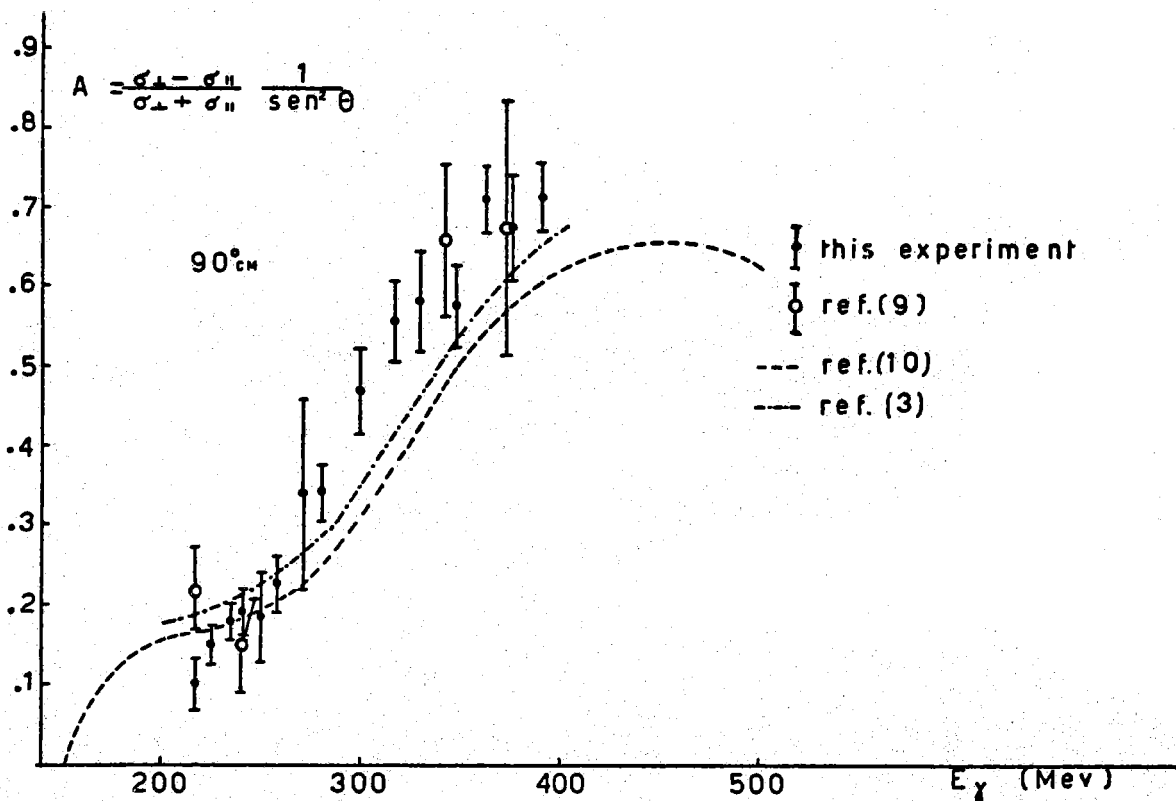


FIG. 2 - Final results for the asymmetry ratio $A(90^\circ)$ at different γ -ray's energies (E_γ). We report also some previous results of Smith and Mozley⁽⁹⁾. For comparison we have indicated the results of the theoretical calculations of Schmidt⁽¹⁰⁾ and of Donnachie and Shaw⁽³⁾.

TABLE I - Asymmetry ratio A (90°)

E_γ (MeV)	A	Reference
217	$.100 \pm .033$	(x)
225	$.150 \pm .025$	(x)
227	$.219 \pm .048$	(o)
235	$.179 \pm .023$	(x)
240	$.191 \pm .029$	(x)
240	$.148 \pm .038$	(o)
250	$.184 \pm .056$	(x)
260	$.226 \pm .035$	(x)
270	$.328 \pm .120$	(x)
281	$.341 \pm .035$	(x)
300	$.469 \pm .054$	(x)
315	$.556 \pm .050$	(x)
330	$.581 \pm .062$	(x)
342	$.664 \pm .098$	(o)
348	$.575 \pm .051$	(x)
363	$.710 \pm .043$	(x)
373	$.671 \pm .164$	(o)
376	$.675 \pm .067$	(x)
391	$.714 \pm .044$	(x)

(x) Present work

(o) Stanford (see (3)).

A complete discussion of our results and of the comparison of these with the theoretical predictions will be done later on the basis of the data relative not only to $\theta = 90^\circ$ but also to different angles and energies. These measurements and the corresponding analysis is now in progress.

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